

3. JEDNADŽBA GIBANJA

SLYJEDI IZ 2. NEWTONOVOG ZAKONA:

SUMA SVIH SILA = BRZINA PROMJENE KOLIČINE
GIBANJA
(INERCIJALNA SILA, $m \cdot a$)

SUMA SVIH MOMENATA = BRZINA PROMJENE
KUĆNE KOLIČINE GIBANJA

SILÉ (VANJSKE, ZADANE)

- OPISUJU DJELOVANJE VANJSKOG SUJETA NA TIJELO

- VANJSKA VOLUMNA SILA:

$$f^p: \Omega^p \times \mathbb{R} \rightarrow \mathbb{R}^3 \quad (\text{ZADRIMO } f^p: \mathbb{R} \rightarrow \mathbb{R}^3)$$

$$\text{ZA } P \subseteq \Omega^p \quad \int_P f^p dx^p \quad - \text{ UKUPNA VOLUMNA SILA NA KOMAD } P$$

f^p - VOLUMNA GUSTOĆA

PRIMJER: SILA TEŽE

- VANJSKA KONTAKTNA SILA:

$$P_1^p \subseteq \partial \Omega^p$$

$$g^p: P_1^p \times \mathbb{R} \rightarrow \mathbb{R} \quad (\text{PLOŠNA}) \text{ GUSTOĆA VANJSKE KONTAKTNE SILÉ}$$

$$\text{ZA } A \subseteq P_1^p \quad \int_A g^p da^p \quad - \text{ UKUPNA VANJSKA KONTAKTNA SILA NA KOMAD } A$$

$\gamma: D \rightarrow A$ PARAMETRIZACIJA

$$\int_A g^p da^p = \int_D g^p \circ \gamma \sqrt{\det(\nabla \gamma^T \nabla \gamma)} dy$$

PRIMER:



- TLAK (NA BALON U VODI)
HPR.

GUSTOĆA MASE:

- $g^p: \Omega^p + \mathbb{R} \rightarrow \mathbb{R}$ NA DEFORMIRANOJ KONFIGURACIJI

- $\varphi \subseteq \Omega^p$ $\int_P g^p da^p$ - MASA KOMPARA P

HAP:

$b^p: \Omega^p + \mathbb{R} \rightarrow \mathbb{R}$ DEF SA $b^p = \frac{f^p}{g^p}$

GUSTOĆA VANJSKE SILE PO JEDINICI MASE

KOLIČINE GIBANJA:

- LINEARNA, DANA GUSTOĆOM $g^p v^p$
(LINEAR MOMENTUM)

- KOTNA, DANA GUSTOĆOM $x^p \times g^p v^p$
(ANGULAR MOMENTUM)

AKSIOM (EULER, CAUCHY)

↙ NAPREŽENJE

POSTOJI VEKTORSKO POLJE $t^p: \overline{\Omega^p} \times \mathbb{R} \rightarrow \mathbb{R}^3$,

$S^2 = \{x \in \mathbb{R}^3: \|x\|=1\}$, TAKVO DA JE

a) $t^p(x^p, u^p, t) = g^p(x^p, t)$, $x^p \in \Gamma_1^p$

u^p - JEDINIČNA NORMALA NA Γ_1^p , VANJSKA STRANA Ω^p

b) ZAKON OČUVANJA IMPULSA:

$\forall \mathcal{P}^p \subseteq \overline{\Omega^p}$ OTVOREN I POUZAN

$$\frac{d}{dt} \int_{\mathcal{P}^p} g^p(x^p, t) v^p(x^p, t) dx^p = \int_{\partial \mathcal{P}^p} t^p(x^p, u^p, t) da^p + \int_{\mathcal{P}^p} f^p(x^p, t) dx^p$$

u^p - JEDINIČNA NORMALA NA $\partial \mathcal{P}^p$, VANJSKA NA \mathcal{P}^p

c) ZAKON OČUVANJA MOMENTA IMPULSA

$\forall \mathcal{P}^p \subseteq \overline{\Omega^p}$ OTVOREN I POUZAN

$$\frac{d}{dt} \int_{\mathcal{P}^p} x^p \times g^p(x^p, t) v^p(x^p, t) dx^p = \int_{\partial \mathcal{P}^p} x^p \times t^p(x^p, u^p, t) da^p + \int_{\mathcal{P}^p} x^p \times f^p(x^p, t) dx^p$$

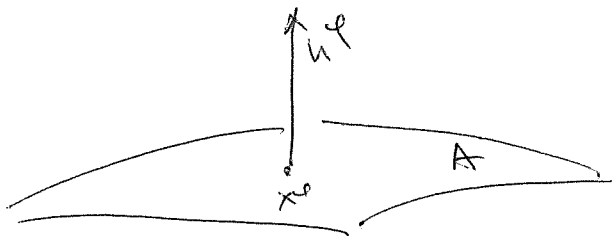
HAP: - t^p SE ZOVE KONTAKTNA SILA

CAUCHYJEV VEKTOR NAPREŽANJA

- ZAPRAVO JE GUSTOĆA

$\int_A t^p da^p$ UKUPNA KONTAKTNA
SILA NA PLOHU A

||
P(A)



$P(A)$ - UKUPNA SILA KOJIM TIJELO
 "IZNAD" (U SMJERU NORMALE n^p)
 DJELUJE NA TIJELO "ISPOD"
 PLOHE A

$$t^p(x^p, n^p, t) \approx \frac{P(A)}{\text{pov}(A)} \rightarrow 0$$

- $\int_{\partial P^p} t^p + \int_{P^p} f^p$ - UKUPNA SILA NA KOMAD P^p

- $\int_{\partial P^p} x^p \otimes t^p + \int_{P^p} x^p \otimes f^p$ - UKUPNI MOMENT NA KOMAD P^p

- KONCEPT KONTAKTNE SILE POSLJEDICA JE FIZIKALNE
 ČIHTJENICE DA JE INTERAKCIJA RAZLIČITIH DJELOVA
 TIJELA POSLJEDICA MEĐUMOLEKULARNIH SILA KOJE
 SU KRATKOG DOHETA

- POSTOJI RIGOROZNI NAČIN DOKAZANJA DA t^p OVISI O
 TOČKI x^p O GEOMETRIJI SAMO KROZ n^p !

- AKSIOM!

ZAKON OČUVANJA MASE

U REFERENTNOJ KONFIGURACIJI $\varphi: \Omega \rightarrow \mathbb{R}^3$

U DEFORMIRANOJ KONFIGURACIJI Ω^t , $\varphi^t: \Omega^t \rightarrow \mathbb{R}^3$

ZAKON OČUVANJA

$$\forall P \subset \Omega \quad \text{ov.} \quad \int_{P^t} \rho^t(x^t, t) dx^t = \int_P \rho(x) dx$$

Pov.

$$P^t = \varphi^t(P) \quad \longrightarrow \quad \int_P \rho^t(\varphi(x, t), t) \det \nabla \varphi(x, t) dx$$

$$\Rightarrow \forall P \quad \int_P \left(\rho(x) - \rho^t(\varphi(x, t), t) \det \nabla \varphi(x, t) \right) dx = 0$$

$$\Rightarrow \quad \rho(x) - \rho^t(\varphi(x, t), t) \det \nabla \varphi(x, t) \quad \left| \frac{\partial}{\partial t} \right.$$

$$\Rightarrow \quad 0 = \frac{\partial}{\partial t} \left(\rho^t(\varphi(x, t), t) \det \nabla \varphi(x, t) \right)$$

KORISTIMO:

$$\frac{\partial}{\partial t} \left(\det \nabla \varphi(x, t) \right) = \left(\operatorname{div} v^t \right) \circ \varphi(x, t) \det \nabla \varphi(x, t)$$

$$\Rightarrow 0 = \frac{\partial}{\partial t} \left(\rho^t(\varphi(x, t), t) \right) \det \nabla \varphi(x, t) + \rho^t(\varphi(x, t), t) \left(\operatorname{div} v^t \right) \circ \varphi(x, t) \det \nabla \varphi(x, t)$$

$$0 = \left(\dot{\rho}^t + \rho^t \operatorname{div} v^t \right) \circ \varphi(x, t) \det \nabla \varphi(x, t)$$

$$\Rightarrow 0 = \dot{\rho}^t + \rho^t \operatorname{div} v^t \quad \text{JEDNAŽBKA KONTINUITETA}$$

TEOREM ZA DOKUJHO GLATKO Q(BANJE P VRJED)

$$\frac{d}{dt} (\det \nabla \varphi(x,t)) = (d\omega v^e) \circ \varphi(x,t) \det \nabla \varphi(x,t)$$

LEMA: $D(\det)(A)(T) = \det A \operatorname{tr}(A^{-1}T)$, $A \in GL(\mathbb{R}^n)$

DOK: \det KAO POLINOM JE DFB.

$$\det X = \det(AA^{-1}X) = \det A \det(A^{-1}X)$$

$$\Rightarrow D(\det)(A)(T) = \det A D(\det)(I)(A^{-1}T)$$

IV: $D(\det)(I) = \operatorname{tr}$

DOK: KORISTENJE DERICIJA U SIJERU

$$D(\det)(I)(T) = \lim_{\varepsilon \rightarrow 0} \frac{\det(I + \varepsilon T) - \det I}{\varepsilon}$$

$$= \lim_{\varepsilon \rightarrow 0} \frac{1 + \varepsilon \operatorname{tr} T + O(\varepsilon^2) - 1}{\varepsilon} = \operatorname{tr} T //$$

$$\Rightarrow D(\det)(A)(T) = \det A \operatorname{tr}(A^{-1}T) //$$

DOKAZ TEOREMA: PRIMENOM LEHE DOBIVAMO

$$\begin{aligned} \frac{d}{dt} (\det \nabla \varphi(x,t)) &= D(\det)(\nabla \varphi(x,t)) \left(\frac{d}{dt} \nabla \varphi(x,t) \right) \\ &= \det \nabla \varphi(x,t) \operatorname{tr} \left(\nabla \varphi_t(x)^{-1} \frac{d}{dt} \nabla \varphi(x,t) \right) \end{aligned}$$

S DRUGE STRANE $\nabla \varphi(x^e, t) = \frac{\partial \varphi}{\partial t}(\varphi_t^{-1}(x^e), t)$

~~$\nabla \varphi(x^e, t) = \nabla \varphi(\varphi_t^{-1}(x^e), t) \nabla \varphi_t^{-1}(x^e)$~~

$$\Rightarrow \nabla \varphi(x^e, t) = \frac{\partial \varphi}{\partial t}(\varphi_t^{-1}(x^e), t) \nabla \varphi_t^{-1}(x^e)$$

$$\Rightarrow d\omega v^e(x^e, t) = \operatorname{tr} \left(\frac{\partial}{\partial t} \nabla \varphi(\varphi_t^{-1}(x^e), t) \nabla \varphi_t^{-1}(x^e) \right) = \operatorname{tr} \left(\nabla \varphi_t^{-1}(x^e) \frac{\partial}{\partial t} \nabla \varphi(\varphi_t^{-1}(x^e), t) \right)$$

$$\& \nabla \varphi_t \circ \varphi_t^{-1} \circ \nabla \varphi_t^{-1} = I //$$

CILJ NAM JE IZVESTI DIFERENCIJALNU FORMULACIJU ZA JEDNAŽBU GIBANJA

LEMA: ZA SKALARNO ILI VEKTORSKO POLJE
 POLE f : I $P \subseteq \mathbb{R}^d$ VRIJEDI TRANSPORTNA
 FORMULA ($\Phi_t^P = \varphi_t(P)$)

$$\frac{d}{dt} \int_{\Phi_t^P} f(x^p, t) dx^p = \int_{\Phi_t^P} \left(\dot{f}(x^p, t) + f(x^p, t) \operatorname{div} v^p(x^p, t) \right) dx^p$$

DOK: ZAMJENOM VARIJABLI $x^p = \varphi_t(x)$
 INTEGRAL OVISAN O t PREBACIMO NA DOMENU P
 NEOVISNU O t .

$$\begin{aligned} \frac{d}{dt} \int_{\varphi_t(P)} f(x^p, t) dx^p &= \frac{d}{dt} \int_P f(\varphi(x, t), t) \det \nabla \varphi(x, t) dx \\ &= \int_P \frac{\partial}{\partial t} \left(f(\varphi(x, t), t) \right) \det \nabla \varphi(x, t) + f(\varphi(x, t), t) \frac{\partial}{\partial t} (\det \nabla \varphi(x, t)) dx \end{aligned}$$

koristimo

$$\frac{\partial}{\partial t} (\det \nabla \varphi(x, t)) = (\operatorname{div} v^p) \circ \varphi(x, t) \det \nabla \varphi(x, t)$$

$$= \int_P \left(\frac{\partial}{\partial t} (f(\varphi(x, t), t)) + f(\varphi(x, t), t) (\operatorname{div} v^p) \circ \varphi(x, t) \right) \det \nabla \varphi(x, t) dx$$

$$= \int_{\Phi_t^P} \left(\dot{f}(x^p, t) + f(x^p, t) \operatorname{div} v^p(x^p, t) \right) dx^p$$

KOROLAR $\frac{d}{dt} \int_{\mathcal{P}^t} s^p v^p dx^p = \int_{\mathcal{P}^t} s^p \dot{v}^p dx^p, \quad \mathcal{P} \subset \Omega, \mathcal{P}^t = \varphi_t(\mathcal{P})$

DOK: PRIMJENOM TRANSPORTNE FORMULE DOBIVAMO

$$\frac{d}{dt} \int_{\mathcal{P}^t} s^p v^p dx^p = \int_{\mathcal{P}^t} (s^p v^p)^\circ + s^p v^p \operatorname{div} v^p dx^p$$

ZA EULEROVA POJA f, g VRIJEDI

$$\begin{aligned} (fg)^\circ(x^p, t) &= \frac{d}{dt} (f(\varphi_t(x), t) g(\varphi_t(x), t)) \circ \varphi_t^{-1}(x^p) \\ &= \frac{d}{dt} (f(\varphi_t(x), t)) \circ \varphi_t^{-1}(x^p) g(x^p, t) \\ &\quad + f(x^p, t) \frac{d}{dt} (g(\varphi_t(x), t)) \circ \varphi_t^{-1}(x^p) = \dot{f} g + f \dot{g} \end{aligned}$$

STOJA VRIJEDI

$$\frac{d}{dt} \int_{\mathcal{P}^t} s^p v^p dx^p = \int_{\mathcal{P}^t} s^p \dot{v}^p + \underbrace{\dot{f}^i v^p + s^p v^p \operatorname{div} v^p}_{=0} dx$$

IZ JEDNAKOSTI

KONTINUITETA

TEOREM (CAUCHY):

HEKA SU $t^p: \Omega^p \times S^2 \times \mathbb{R} \rightarrow \mathbb{R}^3$

$f^p: \Omega^p \times \mathbb{R} \rightarrow \mathbb{R}^3$ HEPREKIDHE

$g^p, v^p: \Omega^p \times \mathbb{R} \rightarrow \mathbb{R}^3$

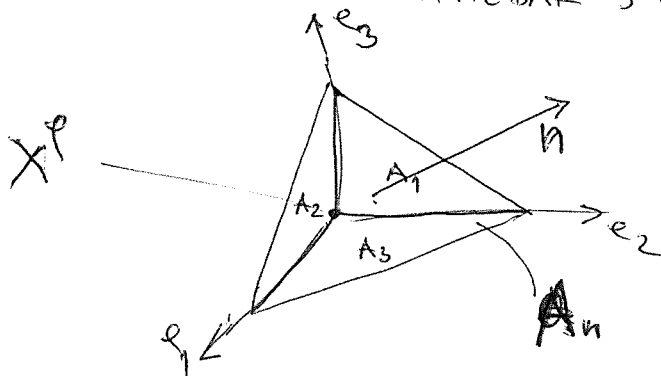
1 HEKA XRYJEDI ZAKON OČUVANJA IMPULSA

TADA $\exists T^p: \Omega^p \times \mathbb{R} \rightarrow \mathbb{H}_3(\mathbb{R})$ T.D
 (CAUCHYJEV TENZOR NAPREZANJA)

$$t^p(x^p, u^p, t) = T^p(x^p, t) n^p, \quad x^p \in \Omega^p, u^p \in S^2$$

DOK: HEKA JE $x^p \in \Omega^p$, $n \in S^2$, $u = u_1 e_1 + u_2 e_2 + u_3 e_3$
 $u_1, u_2, u_3 > 0$

DEF: $A^p =$ TETRAEDAR S VRHOM U x^p KAO NA SLICI



$i \in \{1, 2, 3\}$

TRI STRANICE A_1, A_2, A_3
 U KOORDINATNIM RAVNINAMA
 A_n S NORMALOM n

NORMALA NA $A_j = -e_j, j=1, 2, 3$

$$\int_{\partial A^p} t^p_i(x^p, u^p, t) da^p = \int_{A_1} t^p_i(x^p, -e_1, t) da^p + \int_{A_2} t^p_i(x^p, -e_2, t) da^p + \int_{A_3} t^p_i(x^p, -e_3, t) da^p + \int_{A_n} t^p_i(x^p, n, t) da^p$$

JEK JE t^p NEPREKIDNA MOŽEMO PRIMJENITI TEOREM
SREDNJE VRIJEDNOSTI ZA INTEGRALE, PA POSTOJE TOČKE

$$x^{ij} \in A_j, \quad j=1,2,3, \quad x^i \in A_n$$

TAKVE DA JE

$$\int_{\partial A^p} t^p_i (y^p, u^p, t) da^p = t^p_i(x^{i1}, -e_{1,t}) \text{pov}(A_1) + t^p_i(x^{i2}, -e_{2,t}) \text{pov}(A_2) \\ + t^p_i(x^{i3}, -e_{3,t}) \text{pov}(A_3) + t^p_i(x^i, n, t) \text{pov}(A_n)$$

GEOMETRIJSKI ARGUMENT $\Rightarrow \text{pov}(A_i) = n_i \text{pov} A_n$

$$= (t^p_i(x^{i1}, -e_{1,t}) n_1 + t^p_i(x^{i2}, -e_{2,t}) n_2 + t^p_i(x^{i3}, -e_{3,t}) n_3 + t^p_i(x^i, n, t)) \text{pov} A_n$$

PRIMJENA ZAKONA OČUVANJA IMPULSA SADA POUČAVAMO

$$\left| t^p_i(x^{i1}, -e_{1,t}) n_1 + t^p_i(x^{i2}, -e_{2,t}) n_2 + t^p_i(x^{i3}, -e_{3,t}) n_3 + t^p_i(x^i, n, t) \right|$$

$$\leq \frac{1}{\text{pov} A_n} \left| \int_{A^p} (g^p_{ij} \dot{v}^j - f^p) dx^p \right|$$

$g^p_{ij} \dot{v}^j - f^p$ JE NEPREKIDNA NA $\overline{A^p} \Rightarrow$ OGRANIČENA S C

$$\leq C \frac{\text{vol}(A^p)}{\text{pov}(A_n)} \longrightarrow 0 \quad \text{KAD} \quad A^p \longrightarrow \{x^p\}$$

HOMOTOPNO

S TIH 1

$$x^{ij} \rightarrow x^p$$

$$x^i \rightarrow x^p$$

U LIMESU DOBIVAMO

$$t^p_i(x^p, n, t) = - \sum_{j=1}^3 t^p_i(x^p, -e_j, t) n_j$$

JEZ TO VRIJEDI ZA SVAKU KOMPONENTU

$$t^p(x^p, u, t) = - \sum_{j=1}^3 t^p(x^p, -e_j, t) u_j$$

JEZ JE t^p HEPREKIDNA U LIHESU $u \rightarrow e_k$ ($u_j > 0$)

$$t^p(x^p, e_k, t) = - t^p(x^p, -e_k, t)$$

SADA JE

$$t^p(x^p, u, t) = \sum_{j=1}^3 t^p(x^p, e_j, t) u_j \quad (u_j > 0)$$

DEFINIRAMO:

$$T^p(x^p, t) e_j := t^p(x^p, e_j, t) \quad j=1, 2, 3.$$

SADA VRIJEDI

$$\begin{aligned} T^p(x^p, t) u &= \sum_{j=1}^3 u_j T^p(x^p, t) e_j \\ &= \sum_{j=1}^3 |u_j| T^p(x^p, t) \operatorname{sign}(u_j) e_j \\ &= \sum_{j=1}^3 |u_j| \operatorname{sign}(u_j) t^p(x^p, e_j, t) \\ &= \sum_{j=1}^3 |u_j| t^p(x^p, \operatorname{sign}(u_j) e_j, t) \\ &= \cancel{\sum_{j=1}^3} t^p(x^p, \sum_{j=1}^3 |u_j| \operatorname{sign}(u_j) e_j, t) \\ &= \underline{\underline{t^p(x^p, u, t)}} \end{aligned}$$

ZAKON OČUVANJA IMPULSA SADA GLASI

$$\int_{A^e} g^e \dot{v}^e dx^e = \int_{\partial A^e} T^e(x^e, t) n^e da^e + \int_{A^e} f^e dx^e \quad \forall A^e$$

TEOREM O DIVERGENCIJI OMOGUĆUJE PŘELAZEK NA INTEGRAL PO DOMENI A^e

$$\int_{\partial A} T n da = \int_A dw T dx \quad (\bar{T} - \text{MATICHNÁ FUNKCJA})$$

(DIREKTNÁ POSLYEDICA KLASIČNOLY TEOREMA UZ

$$dw T = \begin{bmatrix} dw & t_1 \\ dw & t_2 \\ dw & t_3 \end{bmatrix}, \quad \bar{T} = \begin{bmatrix} t_1^T \\ t_2^T \\ t_3^T \end{bmatrix}$$

DOBIVAMO

$$\int_{A^e} g^e \dot{v}^e dx^e = \int_{A^e} dw T^e dx^e + \int_{A^e} f^e dx^e$$

$$\int_{A^e} (g^e \dot{v}^e - dw T^e - f^e) dx^e = 0, \quad \forall A^e$$

PROIZVOČNOST DOMENE

$$g^e \dot{v}^e = dw T^e + f^e \quad \text{u } \Omega^e$$

DIFERENCIJALNI OBLIK ZAKONA OČUVANJA IMPULSA

JEDNADŽBA GIBANJA U EULEROVJ FOMULACIJI

POSLEDICE ZAKONA OČUVANJA MOMENTA IMPULSA

ZAKON:

$$\int_{A^e} x^p \times g^p \dot{x}^p dx^p = \int_{\partial A^e} x^p \times T_{ij}^p da^p + \int_{A^e} x^p \times f^p dx^p$$

$$\Rightarrow \int_{A^e} x^p \times (g^p \dot{x}^p - f^p) dx^p = \int_{\partial A^e} x^p \times T_{ij}^p da^p$$

|| ← JEDNAČBA GIBANJA

$$\int_{A^e} x^p \times dw T^p dx^p$$

DESNU STRANU RAČUNAMO KORIŠTEJEM TEOREMA O DIVERGENCIJI

ZA $v \in \mathbb{R}^3 \rightarrow A_v$ - ANTISIMETRIČNA S AKSIJALNIM VEKTOROM v

$$\int_{\partial A^e} x^p \times T^p u^p da^p = \int_{\partial A^e} (A_{x^p} T^p) u^p da^p = \int_{A^e} dw (A_{x^p} T^p) dx^p$$

$$= \int_{A^e} \sum_i \partial_{x_i^p} (x_{x^p} T^p e_i) dx^p = \sum_i \int_{A^e} \partial_{x_i^p} (x_{x^p} \times T^p e_i) dx^p$$

$$= \sum_i \int_{A^e} e_i \times T^p e_i + x^p \times \partial_{x_i^p} (T^p e_i) dx^p$$

$$= \int_{A^e} \sum_i e_i \times T^p e_i + x^p \times dw T^p dx^p$$

GORNJA JEDNADŽBA POVLAGE

$$\int_{A^p} \sum_i e_i \otimes T^p e_i \, dx^p = 0 \quad \forall A^p$$

PROIZVOLJNOST DOMENE POVLAGE

$$\sum e_i \otimes T^p e_i = 0$$

$$e_1 \otimes (T_{11}^p e_1 + T_{21}^p e_2 + T_{31}^p e_3) + e_2 \otimes (T_{12}^p e_1 + T_{22}^p e_2 + T_{32}^p e_3) + e_3 \otimes (T_{13}^p e_1 + T_{23}^p e_2 + T_{33}^p e_3)$$

$$= T_{21}^p e_3 \otimes e_1 + T_{31}^p e_2 \otimes e_1 - T_{12}^p e_3 \otimes e_2 + T_{32}^p e_1 \otimes e_2 + T_{13}^p e_2 \otimes e_3 - T_{23}^p e_1 \otimes e_3 = 0$$

$$\Rightarrow T_{21}^p = T_{12}^p, \quad T_{31}^p = T_{13}^p, \quad T_{32}^p = T_{23}^p$$

$$\Rightarrow T^p \otimes T = T^p$$

TO JE JEDNADŽBA GIBANJA

(CAUCHYJEV TENZOR
NAPREZANJA JE SIMETRIČAN)

POSLEDICE AKSIOMA:

DIFERENCIJALNA JEDNADŽBA GIBANJA U DEFORMIRANOJ KONFIGURACIJI

$$g^p \otimes \dot{v}^p = \operatorname{div} T^p + f^p \quad \text{u } \mathcal{R}^p$$

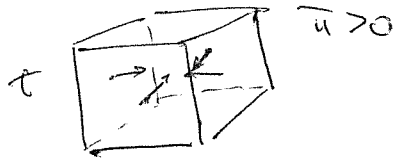
$$T^p \otimes T = T^p \quad \text{u } \mathcal{R}^p$$

$$T^p \otimes u^p = g^p \quad \text{u } \mathcal{P}_1^p$$

PRIMJER (TLAK)

$$T^p(x^e) = -\bar{u} I, \quad \bar{u} \in \mathbb{R} \quad T^p = \begin{bmatrix} -\bar{u} & & \\ & -\bar{u} & \\ & & -\bar{u} \end{bmatrix}$$

$t^p(x^e, n) = -\bar{u} n$ OKOMIT NA PLOHU NA KOJU DJELUJE
ZA $\bar{u} > 0$ U SUPROTHNOJ SMJERU OD n



PRIMJER (ČISTA TENZIJA / KOMPRESIJA)

$$T^p(x^e) = \tau e \otimes e, \quad e \in \mathbb{R}^3, \|e\|=1, \tau \in \mathbb{R}$$

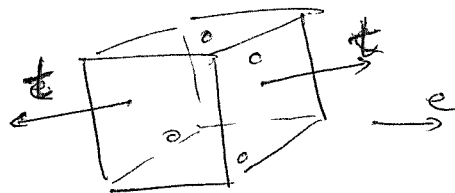
$$t^p(x^e, n) = T^p(x^e)n = \tau(e \cdot n)e$$

- PARALELAN S e

- $\tau > 0$ TENZIJA, $\tau < 0$ KOMPRESIJA

- ZA $n \perp e$ $t^p = 0$

$$T^p = \begin{bmatrix} \tau & & \\ & 0 & \\ & & 0 \end{bmatrix} \quad \text{ZA } e = e_1$$

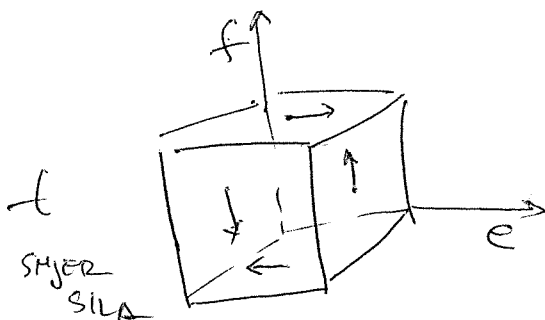


PRIMJER (ČISTO SMICANJE)

$$T^p(x^e, t) = \sigma(e \otimes f + f \otimes e), \quad \sigma \in \mathbb{R}, e, f \in \mathbb{R}^3, \|e\| = \|f\| = 1, e \cdot f = 0$$

$$t^p(x^e, n) = \sigma((f \cdot n)e + (e \cdot n)f)$$

$$T^p = \begin{bmatrix} 0 & \sigma & 0 \\ \sigma & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad \begin{array}{l} e = e_1 \\ f = e_2 \end{array}$$



SLABA FORMULACIJA

$$g^p \dot{x}^p = dw T^p + f^p \quad \text{u } \Omega^p$$

$$T^{pT} = T^p \quad \text{u } \Omega^p$$

$$T^p u^e = g^e \quad \text{u } P_1^p$$

POMHODIMO 1. JABU S θ^p DOVOLJNO GLATKOM I T.D.

$$\theta^p |_{P_0^p} = 0 \quad (P_0^p = \partial \Omega^p - P_1^p)$$

I INTEGRIRAMO PO Ω^p

$$(*) \int_{\Omega^p} g^p \dot{x}^p \cdot \theta^p dx^p = \int_{\Omega^p} dw T^p \cdot \theta^p dx^p + \int_{\Omega^p} f^p \cdot \theta^p dx^p$$

DELIMO PRIMJENITI TEOREM O DIVERGENCIJI

ZA $\eta = T\theta$ RAČUNAMO

$$\begin{aligned} dw \eta &= \partial_1 \eta_1 + \partial_2 \eta_2 + \partial_3 \eta_3 \\ &= \partial_1 (T_{11} \theta_1 + T_{12} \theta_2 + T_{13} \theta_3) + \partial_2 (T_{21} \theta_1 + T_{22} \theta_2 + T_{23} \theta_3) \\ &\quad + \partial_3 (T_{31} \theta_1 + T_{32} \theta_2 + T_{33} \theta_3) \\ &= \partial_1 T_{11} \theta_1 + \partial_1 T_{12} \theta_2 + \partial_1 T_{13} \theta_3 + \partial_2 T_{21} \theta_1 + \partial_2 T_{22} \theta_2 + \partial_2 T_{23} \theta_3 \\ &\quad + \partial_3 T_{31} \theta_1 + \partial_3 T_{32} \theta_2 + \partial_3 T_{33} \theta_3 \\ &\quad + T_{11} \partial_1 \theta_1 + T_{12} \partial_1 \theta_2 + T_{13} \partial_1 \theta_3 + T_{21} \partial_2 \theta_1 + T_{22} \partial_2 \theta_2 + T_{23} \partial_2 \theta_3 \\ &\quad + T_{31} \partial_3 \theta_1 + T_{32} \partial_3 \theta_2 + T_{33} \partial_3 \theta_3 \\ &= \begin{bmatrix} \partial_1 T_{11} + \partial_2 T_{21} + \partial_3 T_{31} \\ \partial_1 T_{12} + \partial_2 T_{22} + \partial_3 T_{32} \\ \partial_1 T_{13} + \partial_2 T_{23} + \partial_3 T_{33} \end{bmatrix} \cdot \theta + T \cdot \begin{bmatrix} \partial_1 \theta_1 & \partial_1 \theta_2 & \partial_1 \theta_3 \\ \partial_2 \theta_1 & \partial_2 \theta_2 & \partial_2 \theta_3 \\ \partial_3 \theta_1 & \partial_3 \theta_2 & \partial_3 \theta_3 \end{bmatrix} \end{aligned}$$

$$\Rightarrow dw(T\theta) = (dw T^T) \cdot \theta + T^T \cdot \nabla \theta$$

ODNOSNO

$$(*) \quad dw(T^T \theta) = (dw T) \cdot \theta + T \cdot \nabla \theta$$

SADA SE VRAĆAMO U (*):

$$\int_{\Omega^e} g^e \dot{v}^e \cdot \theta^e dx^e = \int_{\Omega^e} \text{div}(T^{eT} \theta^e) - T^e \cdot \nabla \theta^e dx^e + \int_{\Omega^e} f^e \cdot \theta^e dx^e$$

(TEOREM O DIVERGENCIJI) = $\int_{\partial \Omega^e} T^{eT} \theta^e \cdot n^e da^e - \int_{\Omega^e} T^e \cdot \nabla \theta^e dx^e + \int_{\Omega^e} f^e \cdot \theta^e dx^e$

||

$\Gamma_0^e \cup \Gamma_1^e$

$$\int_{\partial \Omega^e} \theta^e \cdot T^e n^e da^e = \int_{\Gamma_1^e} \theta^e \cdot T^e n^e da^e = \int_{\Gamma_1^e} \theta^e \cdot g^e da^e$$

↑ Γ_1^e

↑ Γ_1^e

$\forall \theta^e \text{ t.d. } \theta^e|_{\Gamma_0^e} = 0$

$\forall \theta^e \text{ t.d. } T^e n^e|_{\Gamma_1^e} = g^e$

STOGA DOBIVAMO

$$\int_{\Omega^e} g^e \dot{v}^e \cdot \theta^e dx^e + \int_{\Omega^e} T^e \cdot \nabla \theta^e dx^e = \int_{\Omega^e} f^e \cdot \theta^e dx^e + \int_{\Gamma_1^e} g^e \cdot \theta^e da^e$$

(+)

$\forall \theta^e \text{ t.d. } \theta^e|_{\Gamma_0^e} = 0$

VARIJACIJSKA (SLABA) FORMULACIJA JEDNAČBE GIBANJA
U DEFORMIRANOJ KONFIGURACIJI

II PRINCIP VIRTUALNOG RADA U DEFORMIRANOJ KONFIGURACIJI

NAZ: KAO I OBICNO IZ SLABE FORMULACIJE

MOZEMO SE VRATITI U DIFERENCIJALNU

HEXA JE $\theta^p \Big|_{\partial \Omega^p} = 0$. IZ (X) I (T) SLJEDI

$$\int_{\Omega^p} g^p \dot{v}^p \cdot \theta^p dx^p + \int_{\Omega^p} \left(\text{div}(\tau^{pT} \theta^p) - \text{div} \tau^p \cdot \theta^p \right) dx^p = \int_{\Omega^p} f^p \cdot \theta^p dx^p$$

||
0 IZ TEOREMA O DIVERGENCIJI

$$\Rightarrow \int_{\Omega^p} \left(g^p \dot{v}^p - \text{div} \tau^p - f^p \right) \cdot \theta^p dx^p = 0$$

PROIZVOLJNOST OD $\theta^p \Rightarrow$ DIF. J.

$$g^p \dot{v}^p = \text{div} \tau^p + f^p$$

SADA UZEMO $\theta^p \Big|_{\Gamma^p} = 0$, PA IZ (X) I (T) DOBIVAMO

$$\int_{\Omega^p} g^p \dot{v}^p \cdot \theta^p dx^p + \int_{\Omega^p} \left(\text{div}(\tau^{pT} \theta^p) - \text{div} \tau^p \cdot \theta^p \right) dx^p$$

$$= \int_{\Omega^p} f^p \cdot \theta^p dx^p + \int_{\Gamma^p} g^p \cdot \theta^p da^p$$

I DIF. J.

$$\int_{\Omega^p} \text{div}(\tau^{pT} \theta^p) dx^p = \int_{\Gamma^p} g^p \cdot \theta^p da^p$$

$$\Rightarrow \tau_{n^p}^p = g^p$$

$$\int_{\partial \Omega^p} \theta^p \cdot \tau_{n^p}^p da^p = \int_{\Gamma^p} \theta^p \cdot \tau_{n^p}^p da^p$$

TEOREM : RUBNA ZADACIA

$$g^1 \dot{v}^p = dw T^p + f^p \quad u \mathbb{R}^p$$

$$T_{u^p}^p = \bar{g}^p \quad u \mathbb{P}_1^p$$

FORMALNO JE EKUIVALENTNA SLABOJ FORMULACIJI (+).

HAP: ZAKON GČUVANJA IMPULSA (VARIJABILNOST U DOČENI)

PRINCIP VIRTUALNOG RADA (VARIJABILNOST U TEST FUNKCIJI)

U NAŠEM KONTAKTU

IZMEDU JE DIFERENCIJALNA FORMULACIJA,

KOJA ZAHTIJEVA VEĆU REGULARNOST FUNKCIJA,

HO MOŽE SE ZAOPČI!