

Discontinuous Galerkin and Nonconforming in Time Optimized Schwarz Waveform Relaxation for Coupling Heterogeneous Problems

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- 1 Motivation: Application for nuclear waste disposal
- 2 Subdomain time stepping with nonconforming time grids
- 3 Numerical results
- 4 Conclusions

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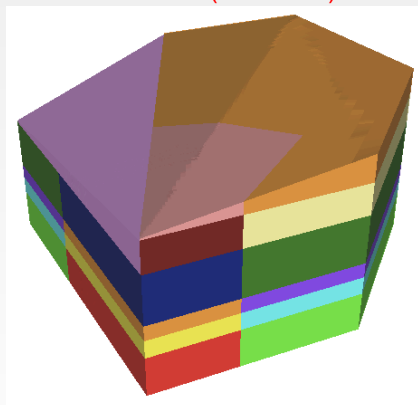
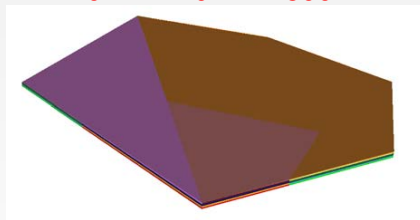
Far field 3D : The computational domain

Research Group MOMAS

with Jérôme Jaffré, Michel Kern and Jean Roberts (INRIA)

A blow-up in the vertical direction (30 times)

Actual dimensions:
 $40\text{km} \times 40\text{km} \times 500\text{m}$



The repository is located in the red part of the bottom layer.

Hydrogeological data

Hydrogeologic layers	Thickness [m]	Porosity [%]	Permeability [m/s]		Effective diffusion coefficient [m ² /s]	Dispersivity Coefficients [m]
			Regional	Local		
Tithonian	Variable	10	$3 \cdot 10^{-5}$	$3 \cdot 10^{-5}$	10^{-9}	6.0, 0.6
Kimmeridgian when it outcrops	Variable	10	$3 \cdot 10^{-4}$	$3 \cdot 10^{-4}$	10^{-9}	6.0, 0.6
Kimmeridgian under cover			10^{-11}	10^{-12}		
Oxfordian L2a-L2b	165	6	$2 \cdot 10^{-7}$	10^{-9}	10^{-9}	6.0, 0.6
Oxfordian Hp1-Hp4	50	18	$6 \cdot 10^{-7}$	$8 \cdot 10^{-9}$	10^{-9}	1600, 30
Oxfordian C3a-C3b	60	1	10^{-10}	10^{-12}	$4 \cdot 10^{-12}$	6.0, 0.6
Callovo-Oxfordian Cox	135	1	$K_v=10^{-14} K_h=10^{-12}$		$4 \cdot 10^{-12}$	6.0, 0.6

⇒ **use different time and space steps, adapted to the physics**

Goal :

- decompose the time interval into windows
- in each window:
 - use an **Optimized Schwarz Waveform Relaxation method**
 - with **non conforming space-time grids**
 - and **discontinuous Galerkin method in time** as subdomain solver

Outline

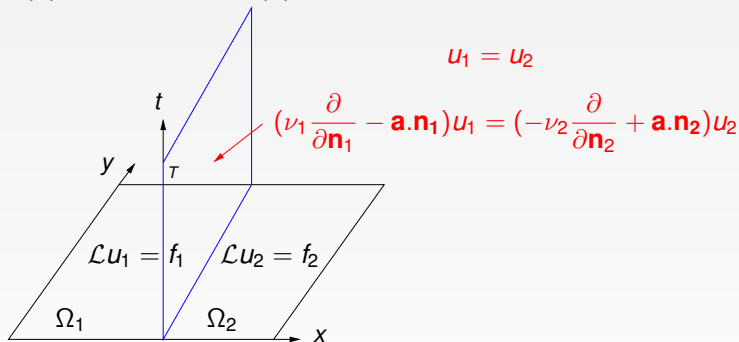
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Advection-diffusion equation with discontinuous coefficients

$$\mathcal{L}u = \frac{\partial u}{\partial t} + bu + \nabla \cdot (\mathbf{a}(\mathbf{x})u - \nu(\mathbf{x})\nabla u) = f \text{ in } \Omega \times [0, T]$$

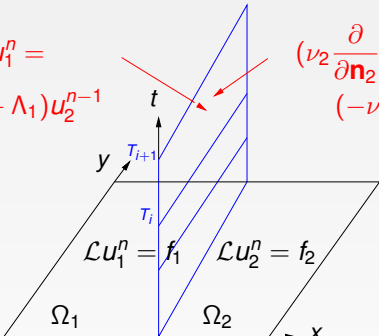
$$u = 0 \text{ on } \partial\Omega \times [0, T], \quad u(\cdot, 0) = u_0, \quad \text{on } \partial\Omega$$

with $\mathbf{a}(\mathbf{x})$ and $\nu(\mathbf{x})$ discontinuous, $\nu(\mathbf{x}) > 0$



Optimized Schwarz Waveform Relaxation Method

(Gander/Halpern/Nataf (DD11, 1998), Martin (2003),
Bennequin/Gander/Halpern (2004), Gander/Halpern/Kern (2004),
Blayo/Halpern/Japhet (2004))

$$\begin{aligned} (\nu_1 \frac{\partial}{\partial \mathbf{n}_1} - \mathbf{a} \cdot \mathbf{n}_1 + \Lambda_1) u_1^n &= \\ (-\nu_2 \frac{\partial}{\partial \mathbf{n}_2} + \mathbf{a} \cdot \mathbf{n}_2 + \Lambda_1) u_2^{n-1} & \end{aligned}$$

$$\begin{aligned} (\nu_2 \frac{\partial}{\partial \mathbf{n}_2} - \mathbf{a} \cdot \mathbf{n}_2 + \Lambda_2) u_2^n &= \\ (-\nu_1 \frac{\partial}{\partial \mathbf{n}_1} + \mathbf{a} \cdot \mathbf{n}_1 + \Lambda_2) u_1^{n-1} & \end{aligned}$$

Choose Λ_1 and Λ_2 in order to optimize the convergence rate

Optimized Schwarz Waveform Relaxation Method

$$\Lambda_1 = \alpha_2 + \beta_2(\partial_t + \mathbf{a} \cdot \boldsymbol{\tau}_2 \partial_{\boldsymbol{\tau}_2} - \partial_{\boldsymbol{\tau}_2}(\nu_2 \partial_{\boldsymbol{\tau}_2})), \quad \Lambda_2 = \alpha_1 + \beta_1(\partial_t + \mathbf{a} \cdot \boldsymbol{\tau}_1 \partial_{\boldsymbol{\tau}_1} - \partial_{\boldsymbol{\tau}_1}(\nu_1 \partial_{\boldsymbol{\tau}_1}))$$

where $\alpha_1, \alpha_2, \beta_1, \beta_2$ optimize the convergence rate

$(\nu_1 \frac{\partial}{\partial \mathbf{n}_1} - \mathbf{a} \cdot \mathbf{n}_1 + \Lambda_1) u_1^n =$
 $(-\nu_2 \frac{\partial}{\partial \mathbf{n}_2} + \mathbf{a} \cdot \mathbf{n}_2 + \Lambda_1) u_2^{n-1}$

$(\nu_2 \frac{\partial}{\partial \mathbf{n}_2} - \mathbf{a} \cdot \mathbf{n}_2 + \Lambda_2) u_2^n =$
 $(-\nu_1 \frac{\partial}{\partial \mathbf{n}_1} + \mathbf{a} \cdot \mathbf{n}_1 + \Lambda_2) u_1^{n-1}$

Ω_1 Ω_2

$\mathcal{L}u_1^n = f_1$ $\mathcal{L}u_2^n = f_2$

τ_i τ_{i+1}

t

y

x

How to discretize these conditions with nonmatching grids in time ?

Discontinuous Galerkin in time

(Eriksson-Johnson-Thomé, 1985, Halpern-Japhet, 2005)

Non conforming finite elements in space

(Gander-Japhet-Maday-Nataf, 2004)

Subdomain problem in Ω_j ,
in one time window $I = (T_i, T_{i+1})$

$$\left\{ \begin{array}{ll} \mathcal{L}u = f & \text{in } \Omega_j \times I, \\ u(\cdot, T_i) = u_0 & \text{in } \Omega_j, \\ \left(\nu \frac{\partial}{\partial \mathbf{n}} - \mathbf{a} \cdot \mathbf{n} + \alpha + \beta \left(\frac{\partial}{\partial t} + \delta \frac{\partial}{\partial \tau} - \frac{\partial}{\partial \tau} \left(r \frac{\partial}{\partial \tau} \right) \right) \right) u = g & \text{on } \Gamma \times I \end{array} \right.$$

Weak Formulation

Let $H_S^s(\Omega) = \{v \in H^s(\Omega), v|_\Gamma \in H^s(\Gamma)\}$, equipped with

$$(u, v)_{H_S^s(\Omega)} = (u, v)_{H^s(\Omega)} + \beta(u, v)_{H^s(\Gamma)}$$

Find u such that

$$(\partial_t u, v)_{H_0^0(\Omega)} + a(u, v) = \ell(v), \quad \forall v \in H_1^1(\Omega)$$

with

$$\begin{cases} a(u, v) = \int_{\Omega} \nabla \cdot (\mathbf{a}u)v \, dx + \int_{\Omega} \nu \nabla u \cdot \nabla v \, dx + \int_{\Omega} buv \, dx \\ \quad + \int_{\Gamma} ((\alpha - \mathbf{a} \cdot \mathbf{n})uv) + \beta \partial_\tau uv + \beta r \partial_\tau u \partial_\tau v \, ds \\ \ell(v) = (f, v)_{L^2(\Omega)} + (g, v)_{L^2(\Gamma)} \end{cases}$$

Time Discontinuous Galerkin

Let \mathcal{T} be a decomposition of $I = \cup_{k=1}^K I^k$ with $I^k = [t_k, t_{k+1}]$. We define

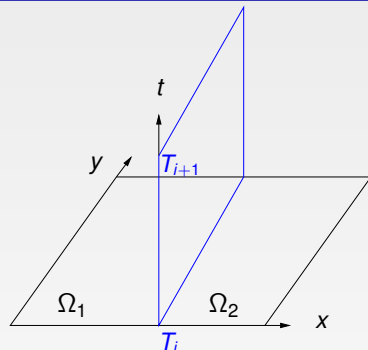
$$\begin{aligned}\mathbf{P}_q(V) &= \left\{ \varphi : \varphi(t) = \sum_{i=0}^q \varphi_i t^i, \varphi_i \in V \right\} \\ \mathcal{P}_q(V, \mathcal{T}) &= \left\{ \varphi : I \rightarrow V, \varphi|_{I^k} \in \mathbf{P}_q(V), 0 \leq k \leq K \right\}.\end{aligned}$$

Let $\varphi(t_k^\pm) = \lim_{t \rightarrow t_k \pm 0} \varphi(t)$

The discontinuous Galerkin method defines recursively on I_k , an approximate solution U in $\mathcal{P}_q(H_1^1(\Omega), \mathcal{T})$ such that

$$\left\{ \begin{array}{l} U(0, \cdot) = u_0, \\ \forall \varphi \in \mathcal{P}_q(H_1^1(\Omega), \mathcal{T}) : \int_{I_k} \left[\left(\frac{dU}{dt}, \varphi \right)_{H_0^0(\Omega)} + a(U, \varphi) \right] dt \\ \quad + \left((U(t_k^+, \cdot) - U(t_k^-, \cdot)), \varphi(t_k^+, \cdot) \right)_{H_0^0(\Omega)} = \int_{I_k} \ell(\varphi) dt \end{array} \right.$$

Domain decomposition



The continuous matching conditions for Ω_1 is

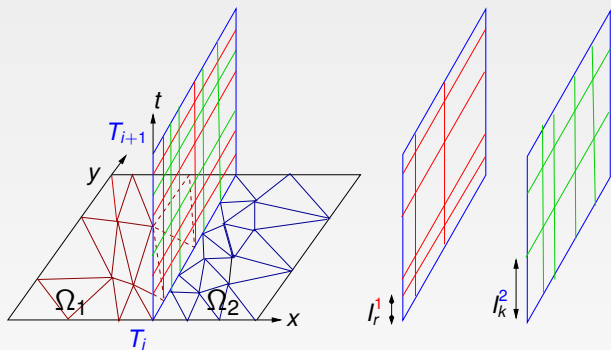
$$f_1(t) = g_2(t), \quad \forall t \in I$$

with :

$$\begin{aligned} f_1(t) &= (\nu_1 \partial_{\mathbf{n}_1} - \mathbf{a} \cdot \mathbf{n}_1 + \Lambda_1) u_1^n \\ g_2(t) &= (-\nu_2 \partial_{\mathbf{n}_2} + \mathbf{a} \cdot \mathbf{n}_2 + \Lambda_1) u_2^{n-1} \end{aligned}$$

Projections between time grids

L^2 orthogonal projection on $\mathcal{P}_q(\mathbb{R}, \mathcal{T}_1)$, restricted to $\mathcal{P}_q(\mathbb{R}, \mathcal{T}_2)$

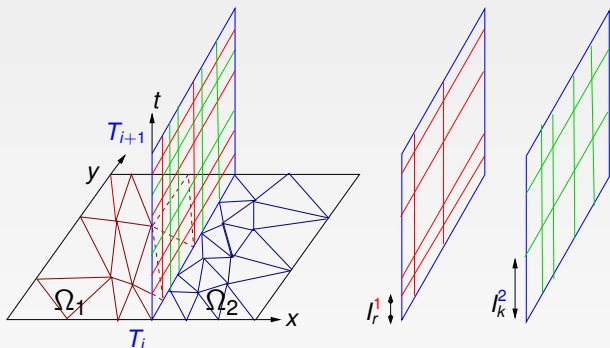


The discrete approximations F_1 of f_1 in $\mathcal{P}_q(\mathbb{R}, \mathcal{T}_1)$, and G_2 of g_2 in $\mathcal{P}_q(\mathbb{R}, \mathcal{T}_2)$ verify the nonconforming matching condition

$$\int_I [F_1 - G_2] V_1 = 0, \quad \forall V_1 \in \mathcal{P}_q(\mathbb{R}, \mathcal{T}_1)$$

An efficient way to perform the projections between time grids

Based on the method in Gander/Halpern/Nataf (2003)



Let V_k^1 (resp. V_ℓ^2) the shape functions of $\mathcal{P}_q(\mathbb{R}, T_1)$ (resp. $\mathcal{P}_q(\mathbb{R}, T_2)$)

How to compute $M_{k,\ell} = \int_I V_k^1 V_\ell^2$?

\Rightarrow Linear complexity algorithm without an additional grid

Convergence - Error estimates

based on the theoretical results in [Eriksson/Johnson/Larsson \(1998\)](#), [Makridakis/Akrivis \(2004\)](#), [Szeftel \(2004\)](#)

Convergence : the continuous algorithm converges for

- $\beta_1 = \beta_2 = 0$, $\alpha_1 \neq \alpha_2$, $\nu_1 \neq \nu_2$, $\mathbf{a}_1 \neq \mathbf{a}_2$ and general decomposition
- if $\beta_1 \neq 0$, $\beta_2 \neq 0$, $\beta_1 = \beta_2$, $\alpha_1 \neq \alpha_2$, $\nu_1 \neq \nu_2$, $\mathbf{a}_1 \neq \mathbf{a}_2$ and decomposition into strips

The coupled discret problem in time, has a unique solution and the discret Schwarz algorithm is convergent.

Error estimates : for $\beta_1 = \beta_2 = 0$

$$\sum_{i=1}^I \|u - U_i\|_{L^\infty(0,T,L^2(\Omega_i))}^2 = \mathcal{O}(\Delta t^{q+1})$$

with $\Delta t = \sup_k \Delta t_k$

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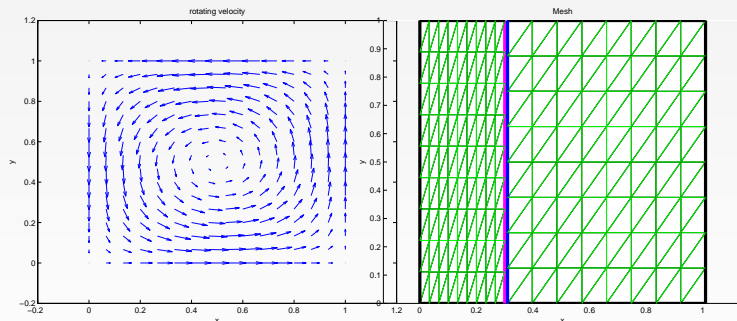
Numerical Results

Replace $H_1^1(\Omega_j)$ with V_j^h (P_1 finite element space)

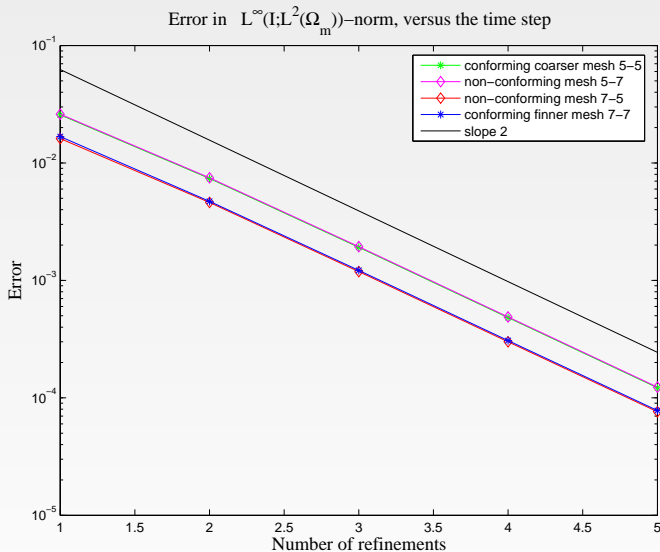
Exact solution $u(x, t) = \cos(\pi x)\sin(\pi y)\cos(\pi t)$, in $[0, 1]^3$

$\mathbf{a} = (-\sin(\pi * (y - \frac{1}{2})). * \cos(\pi * (x - \frac{1}{2})), \cos(\pi * (y - \frac{1}{2})). * \sin(\pi * (x - \frac{1}{2}))),$
 $\nu_1 = \nu_2 = 1$

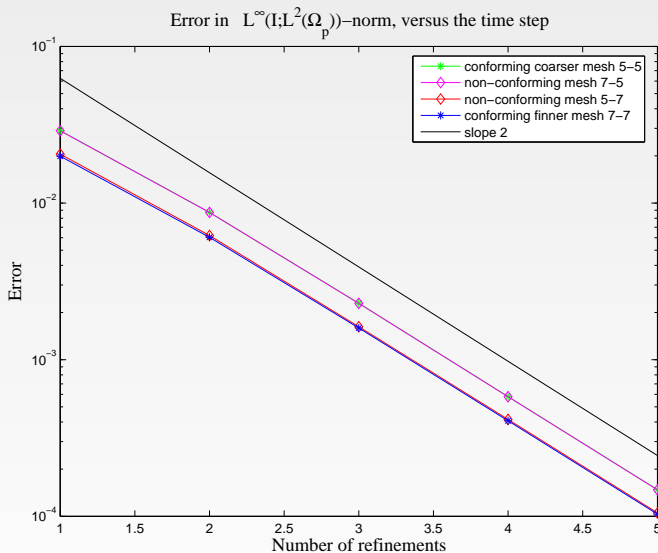
Stopping criterion : the jump of interface conditions is smaller than 10^{-6}
Space-time non conforming grids



Error in $L^\infty(I; L^2(\Omega_1))$ norm



Error in $L^\infty(I; L^2(\Omega_2))$ norm

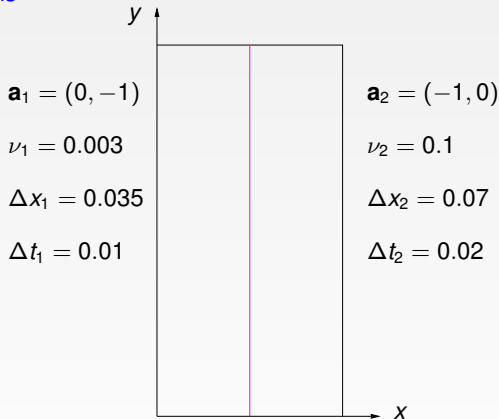


Example with discontinuous coefficients

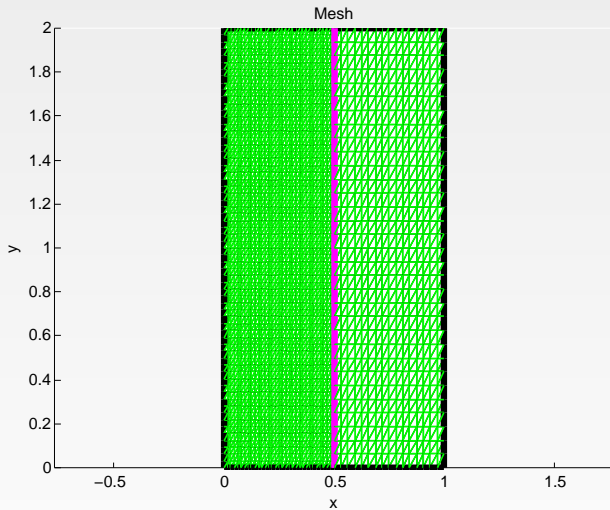
$$\mathcal{L}u = \frac{\partial u}{\partial t} + \nabla \cdot (\mathbf{a}(\mathbf{x})u - \nu \nabla u) = e^{-100((x-0.55)^2 + (y-1.7)^2)} \text{ in } \Omega \times [0, T]$$

$$u = 0 \text{ on } \Gamma_0 \times [0, T], \quad \partial_n u = 0 \text{ on } \partial\Omega \setminus \Gamma_0 \times [0, T], \quad u(\cdot, 0) = e^{-100((x-0.55)^2 + (y-1.7)^2)} \text{ on } \partial\Omega$$

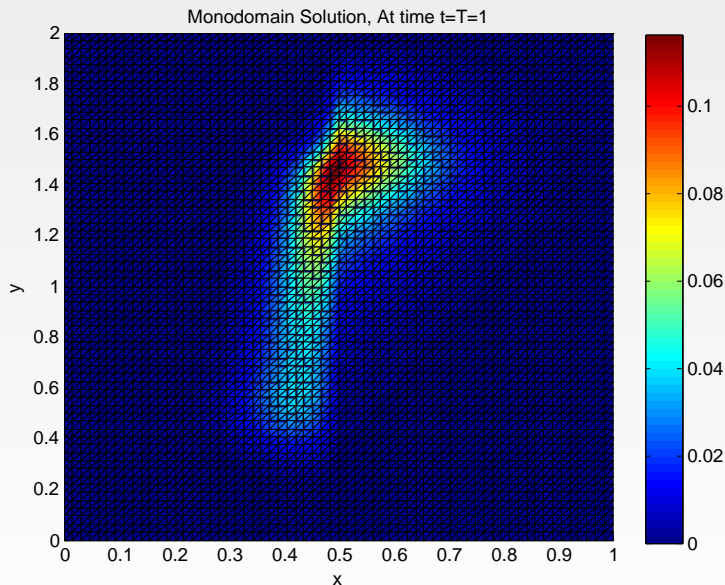
3 OSWR iterations



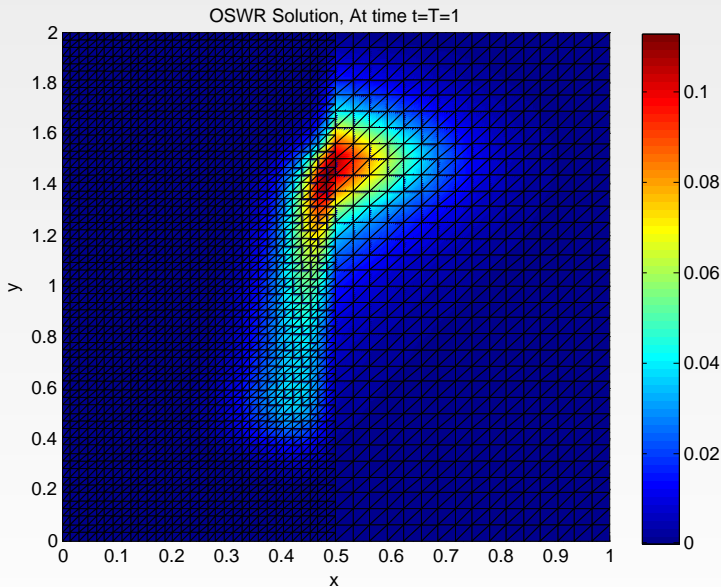
Non conforming space grid



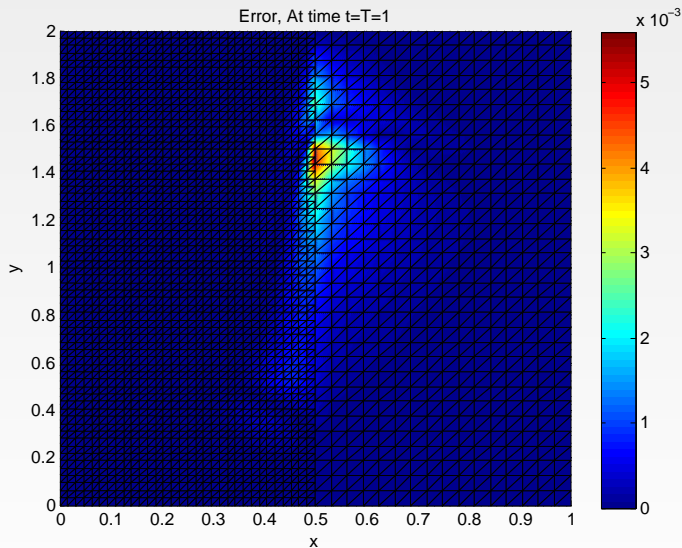
Monodomain Solution at time $T=1$



OSWR Solution at time $T=1$



Error between monodomain and multidomain solutions at time $T=1$



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- Time discontinuous Galerkin method with Optimized Schwarz Waveform Relaxation
 - ⇒ lead to physical transmission conditions in very few iterations
 - ⇒ independant time steps with preservation of the scheme global order in time in the subdomains
 - ⇒ a simple and efficient algorithm to perform projection between nonmatching time grids
- Work in progress
 - numerical and mathematical analysis of the convergence rate (with M.J. Gander)
 - Extension to the MOMAS approach (Mixte Finite Element)